

Matrix hydrodynamics and 2-D Euler

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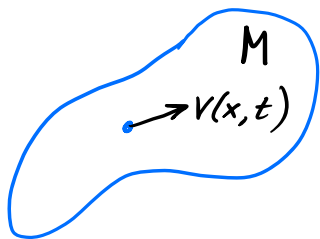


Dedicated to Vladimir Zeitlin

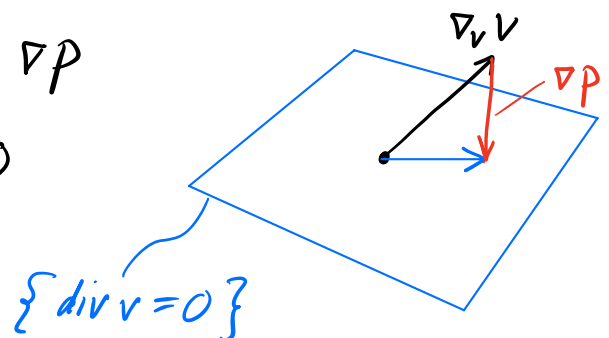


The Father of Matrix Hydrodynamics

Incompressible inviscid hydrodynamics (Euler 1757)



$$\begin{cases} \dot{v} + \nabla_v v = -\nabla p \\ \operatorname{div} v = 0 \end{cases}$$



Arnold's viewpoint

Integrate: $\dot{x}(t) = v(x(t), t)$

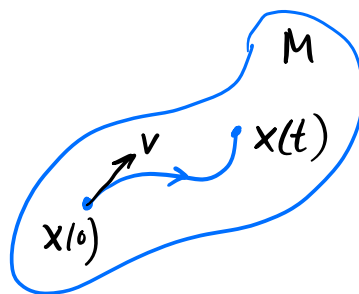
Flow map (all $x \in M$ at once):

$$\dot{\varphi}_t = V_t \circ \varphi_t, \quad \varphi_0(x) = x$$

tracks fluid particles

NOTICE:

$$\varphi_t \in \text{Diff}_\mu(M) = \{\eta \in \text{Diff}(M) \mid \eta_* \mu = \mu\}$$



Newton's equation (differentiate)

$$\ddot{\varphi} = \dot{v} \circ \varphi + \underbrace{\nabla v \circ \varphi}_{v \circ \varphi} \cdot \dot{\varphi} = (\dot{v} + \nabla_v v) \circ \varphi$$

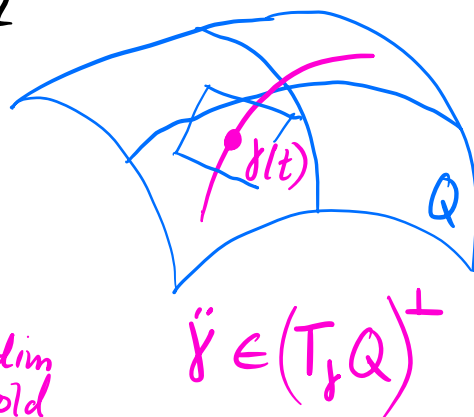
$$= (-\cancel{\nabla_v v} - \nabla p + \cancel{\nabla_v v}) \circ \varphi$$

$$\Rightarrow \ddot{\varphi} = -\nabla p \circ \varphi \quad (*)$$

What is a geodesic curve?

Idea: take $Q = \text{Diff}_\mu(M)$

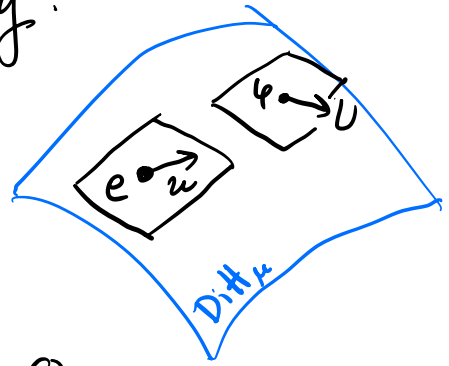
int. dim manifold



Is $\nabla p \circ \varphi \perp$ to $T_\varphi \text{Diff}_\mu \subset T_\varphi \text{Diff}$?

Riemannian metric: kinetic energy:

$$\langle U, U \rangle_\varphi = \int_M |U(x)|^2 d\mu(x)$$



$$U \in T_\varphi \text{Diff}_\mu(M) \iff \text{div}(\underbrace{U \circ \varphi^{-1}}_u) = 0$$

$$U: M \rightarrow TM$$

$$x \mapsto U(x) \in T_{\varphi(x)} M$$

Back to Euler equations: $\ddot{\varphi} = -\nabla P \circ \varphi$

Thm (Arnold 1966)

$\varphi(t) \in \text{Diff}_\mu(M)$ is geodesic curve

proof:

$$\begin{aligned} \langle U, \nabla P \circ \varphi \rangle_\varphi &= \int_M U \cdot \nabla P \circ \varphi d\mu = \int_M (u \cdot \nabla P) \circ \varphi d\mu \\ &= \int_M u \cdot \nabla P \underbrace{\varphi_* d\mu}_{d\mu} = \int_M u \cdot \nabla P d\mu = 0 \quad \square \end{aligned}$$

$$\langle u, \nabla P \rangle_e = \langle u, \nabla P \rangle_{L^2}$$

NOTE: Metric is right-invariant: $\langle U \circ \eta, V \circ \eta \rangle_{\varphi \circ \eta} = \langle U, V \rangle_\varphi$

Hamiltonian description: Lie-Poisson dynamics

G : Lie group

Hamiltonian $\bar{H}: T^*G \rightarrow \mathbb{R}$ is right-invariant:

$$\bar{H}((q,p) \cdot g) = \bar{H}(q,p)$$

Symmetry group G : dynamics reduce to $T^*G/G \simeq \mathfrak{g}^*$

dual of Lie algebra

$$[(q,p)] \longleftrightarrow (q,p) \cdot q^{-1} = (e, \underbrace{p \cdot q^{-1}}_m) \simeq m \in T_e^*G = \mathfrak{g}^*$$

$$\bar{H}(q,p) = H(\underbrace{p \cdot q^{-1}}_m), \quad H: \mathfrak{g}^* \rightarrow \mathbb{R}$$

Hamilton's eq:

$$\begin{cases} \dot{q} = \frac{\partial H}{\partial p} \\ \dot{p} = -\frac{\partial H}{\partial q} \end{cases}$$

$$m \equiv p \cdot q^{-1} \implies$$

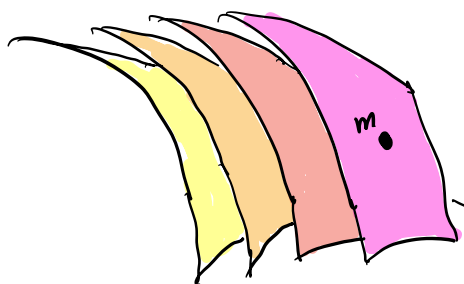
Lie-Poisson system

$$m - \text{ad}_{dH(m)}^* m = 0$$

Preserves Poisson structure on $\mathfrak{g}^*$

$$\{dF, dG\}(m) = \langle m, [dF(m), dG(m)] \rangle$$

$\mathfrak{g}^*$ foliated in co-adjoint orbits (symplectic leaves)



$$O_m = \{ \text{Ad}_g^*(m) \mid g \in G \}$$

Special case M Riemannian (compact)

$$G = \text{Diff}_\mu(M) \quad (\text{Frechet Lie group})$$

$$\mathfrak{g} = \mathfrak{X}_\mu(M) = \{v \in \mathfrak{X}(M) \mid \text{div } v = 0\}$$

$$\mathfrak{g}^* \simeq \Omega^1(M) / d\Omega^0(M) \simeq d\Omega^1(M) \quad (\text{exact 2-forms})$$

smooth
dual

1-trivial
1-co-homology

Hamiltonian on $\mathfrak{g}^* \simeq d\Omega^1(M)$

$$H(\hat{\omega}) = \frac{1}{2} \int_M i_v d^{-1} \hat{\omega} \mu, \quad dV^b = \hat{\omega}, \quad \text{div } v = 0$$

exact
2-form

2-D case:

$$\hat{\omega} = \omega \mu$$

$$v = X_\psi = \nabla^\perp \psi$$

symplectic structure on M

Hamiltonian on M

$$H(\omega) = \frac{1}{2} \int_M \psi \omega \mu, \quad -\Delta \psi = \omega$$

Poisson bracket on M

$$\langle \text{ad}_{X_\psi}^* \hat{\omega}, X_f \rangle = \langle \omega, \{\psi, f\} \rangle_{L^2} = \langle -\{\psi, \omega\}, f \rangle_{L^2}$$

$$\dot{\hat{\omega}} - \text{ad}_{X_\psi}^* \hat{\omega} = 0 \iff$$

$$\dot{\omega} + \{\psi, \omega\} = 0 \quad (*)$$

$$\iff \dot{\omega} + \mathcal{L}_{X_\psi} \omega = 0 \iff \dot{\omega} + \text{div}(\omega \nabla^\top \psi) = 0$$

Casimirs: $C_f(\omega) = \int_M f \circ \omega \mu$ preserved
for each $f: \mathbb{R} \rightarrow \mathbb{R}$

Proof: C_f constant on co-adjoint orbits

Zeitlin's idea (1991)

Use **quantization theory** to replace $(*)$ by
finite dim. Lie-Poisson system

Quantization in a nutshell (for $M = S^2$)

Vorticity state $\omega: S^2 \rightarrow \mathbb{R}$

Aim: mapping $T_N: \omega \mapsto W \in U(N)$ such that

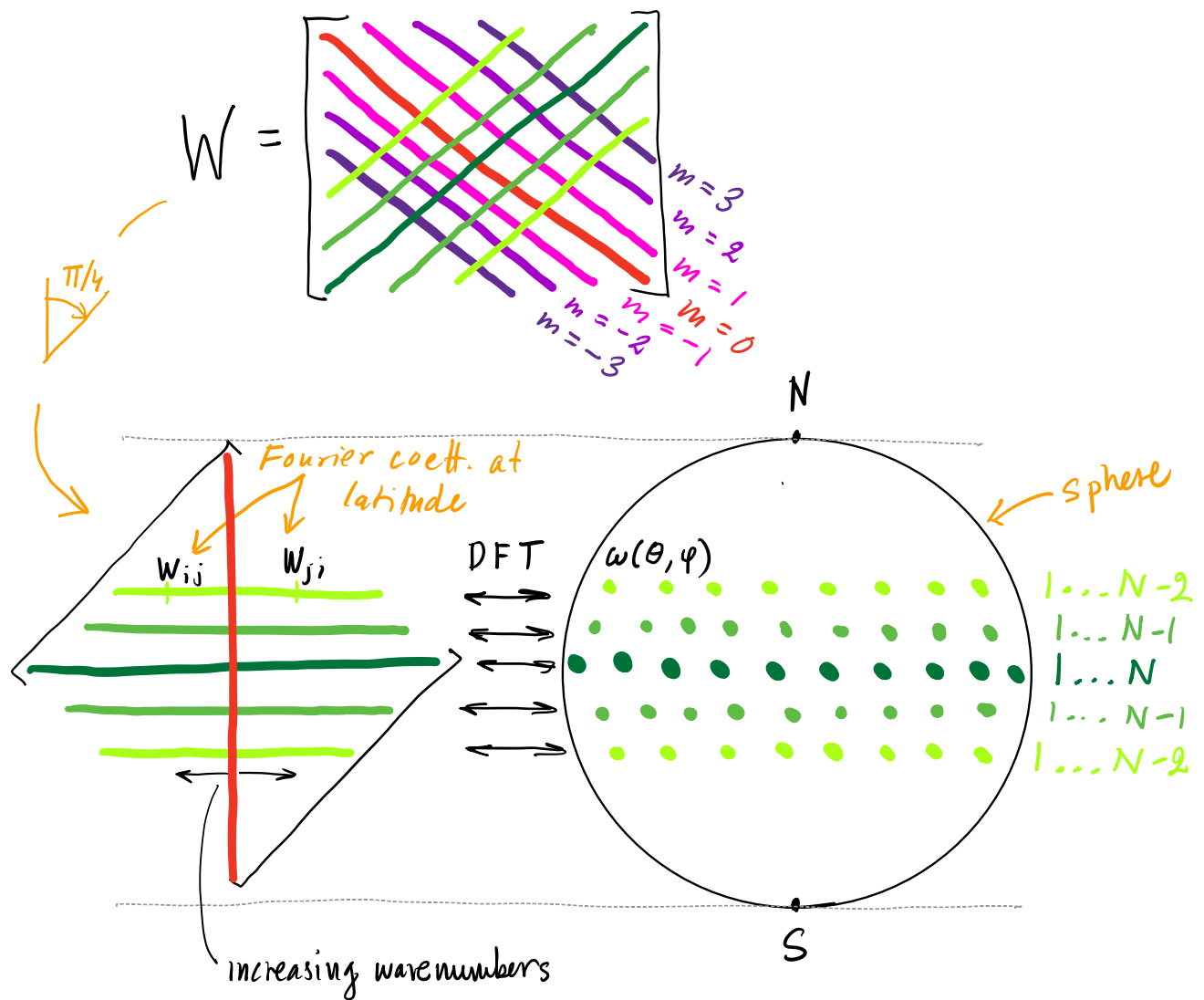
$$T_N(\{\omega_1, \omega_2\}) \approx \frac{1}{\hbar} [T_N \omega_1, T_N \omega_2] \leftarrow \text{matrix commutator}$$

$$\hbar_N = \frac{1}{\sqrt{N^2 - 1}}$$

Hoppe (1989): explicit quantization on π^2 and S^2

Route: $\omega \mapsto \sum_{l=0}^{\infty} \sum_{m=-l}^l \omega_{lm} Y_{lm} \mapsto \sum_{l=0}^N \sum_{m=-l}^l \omega_{lm} T_{lm}^N$
(for S^2)
↑
spherical harmonics
 $\in U(N)$

What is the Matrix? (Neo, 1999)



Euler-Zeitlin equation

$$\omega \in C^\infty(\mathbb{S}^2) \longleftrightarrow W \in \mathcal{U}(N)$$

$$\Delta \psi = \omega \longleftrightarrow \Delta_N P = W$$

Hoppe-Yau Laplacian (1998)

$$\dot{\omega} = \{\psi, \omega\} \longleftrightarrow \dot{W} = \frac{1}{\hbar} [P, W]$$

stream matrix

vorticity matrix

Dictionary Hydrodynamics — Matrix theory

"classical" hydrodynamics

vorticity $w \in L^\infty(\mathbb{S}^2)$

Casimir $I_f(w) = \int_{\mathbb{S}^2} f(w(x)) d\mu(x)$

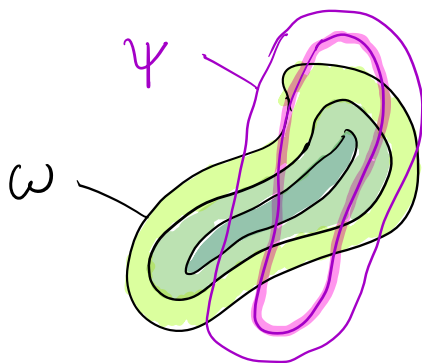
Hamiltonian $H(w) = \frac{1}{2} \int_{\mathbb{S}^2} \psi w d\mu$

values of w

levelsets of w

$\|w\|_{L^\infty}$

average w along levelsets of ψ



matrix hydrodynamics

$W \in \mathfrak{u}(N)$

$I_f^N(W) = \text{tr}(f(W))$

$H^N(W) = \frac{1}{2} \text{tr}(PW)$

eigenvalues of W

eigenvectors of W

spectral norm

projection of W onto

$\text{stab}_P = \{W_s \in \mathfrak{u}(N) \mid [P, W_s] = 0\}$